

Large time behavior of solutions to Schrödinger equations with nonlinearity of integral type

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This talk is based on my recent joint work with Y. Yamazaki [32]. We study the Cauchy problem for the nonlinear Schrödinger equation with interaction described by the integral with respect to one direction of the intensity in two space dimensions:

$$i\partial_t u + \frac{1}{2}\Delta u = f(u), \quad (NLS)$$

where $u : \mathbb{R} \times \mathbb{R}^2 \ni (t, (x, y)) \mapsto u(t, x, y) \in \mathbb{C}$, Δ is the Laplacian in space $\mathbb{R}^2$, and $f(u)$ is the nonlinear interaction given by

$$(f(u))(t, x, y) = \lambda u(t, x, y) \int_{-\infty}^x |u(t, x', y)|^2 dx'$$

with $\lambda \in \mathbb{R}$. For $\alpha, \beta \geq 0$, we define $\dot{H}^{\alpha, \beta} = L_y^2 \dot{H}_x^\alpha \cap L_x^2 \dot{H}_y^\beta$ with norm $\|u; \dot{H}^{\alpha, \beta}\| = \text{Max}(\|u; L_y^2 \dot{H}_x^\alpha\|, \|u; L_x^2 \dot{H}_y^\beta\|)$.

Theorem 1. *Let α and β satisfy $0 \leq \alpha < 1/2$ and $\beta \geq 0$. Let $\phi \in \mathcal{F}(\dot{H}^{\alpha, \beta})$ where $\mathcal{F}$ is the Fourier transform. The (NLS) has a unique solution*

$$u \in X_{\text{loc}} \equiv C(\mathbb{R}; L^2) \cap \bigcap_{0 \leq 2/q = 1/2 - 1/r \leq 1/2} L_{\text{loc}}^q(\mathbb{R}; L_x^2 L_y^r).$$

Moreover, $|J_x|^\alpha u, |J_y|^\beta u \in X_{\text{loc}}$, where

$$\begin{aligned} |J_x|^\alpha &= U_x(t)|x|^\alpha U_x(-t) = M_x(t)(-t^2 \Delta_x)^{\alpha/2} M_x(-t), \\ |J_y|^\beta &= U_y(t)|y|^\beta U_y(-t) = M_y(t)(-t^2 \Delta_y)^{\beta/2} M_y(-t), \\ U_x(t) &= \exp\left(\frac{it}{2} \Delta_x\right), U_y(t) = \exp\left(\frac{it}{2} \Delta_y\right), \\ M_x(t) &= \exp\left(\frac{ix^2}{2t}\right), M_y(t) = \exp\left(\frac{iy^2}{2t}\right). \end{aligned}$$

Basic estimates for the proof of Theorem 1 are based on the fractional Leibniz rule [25], the generalized Hölder inequality in the Lorentz spaces [29,30], and the boundedness of the Hilbert transform and the Riesz potential in the Lorentz space [27,36].

To describe the large time behavior of solutions of (NLS) with small Cauchy data, we introduce modified free dynamics for $\phi_{\pm} \in L^2 \cap \mathcal{F}(L_x^2 L_y^\infty)$:

$$\begin{aligned} v_1^\pm(t) &= U(t) \exp(-iS_\pm(t, -i\nabla))\phi_\pm, \\ v_2^\pm(t) &= U(t)M(-t) \exp(-iS_\pm(t, -i\nabla))\phi_\pm \\ &= M(t)D(t) \exp(-iS_\pm(t, \cdot))\hat{\phi}_\pm, \\ v_3^\pm(t) &= \exp(-iS_\pm(t, t^{-1}x, t^{-1}y))U(t)\phi_\pm \\ &= M(t)D(t) \exp(-iS_\pm(t, \cdot))\mathcal{F}M(t)\phi_\pm, \end{aligned}$$

where

$$\begin{aligned} U(t) &= \exp\left(\frac{it}{2}\Delta\right) = M(t)D(t)\mathcal{F}M(t), \\ M(t) &= \exp\left(\frac{i}{2t}(x^2 + y^2)\right), \\ (D(t)\psi)(x, y) &= (it)^{-1}\psi(t^{-1}x, t^{-1}y), \\ S_\pm(t, x, y) &= \pm\lambda \int_{-\infty}^x |\hat{\phi}_\pm(x', y)|^2 dx' \log |t|. \end{aligned}$$

For $\rho, \rho_0 > 0$ we define

$$B(\rho, \rho_0) = \{\phi \in L^2 \cap \mathcal{F}(\dot{H}^{\alpha, \beta}); \|\phi : L^2\| \leq \rho_0, \|\phi; \mathcal{F}(\dot{H}^{\alpha, \beta})\| \leq \rho\}.$$

Theorem 2. *Let α and β satisfy $0 < \alpha < 1/2 < \beta < 1$. Then for any $\rho_0 > 0$ there exists $\rho > 0$ such that for any $\phi \in B(\rho, \rho_0)$ the solution u given by Theorem 1 satisfies*

$$\begin{aligned} u &\in X \equiv C(\mathbb{R}; L^2) \cap \bigcap_{0 \leq 2/q=1/2-1/r \leq 1/2} L^q(\mathbb{R}; L_x^2 L_y^r), \\ |J_x|^\alpha u, |J_y|^\beta u &\in X, \\ \|u(t); L_x^2 L_y^\infty\| &= O(|t|^{-1/2}) \quad \text{as } t \rightarrow \pm\infty. \end{aligned}$$

Moreover, there exist unique $\phi_\pm \in L^2 \cap \mathcal{F}(L_x^2 L_y^\infty)$ such that for $\varepsilon > 0$ sufficiently small

$$\|\mathcal{F}U(-t)u(t) - \exp(-iS_\pm(t, \cdot))\hat{\phi}_\pm; L^2 \cap L_x^2 L_y^\infty\| = O(|t|^{-\varepsilon})$$

as $t \rightarrow \pm\infty$. Furthermore,

$$\|u(t) - v_j(t); L^2\| = O(|t|^{-\varepsilon})$$

as $t \rightarrow \pm\infty$, for $j = 1, 2$, and

$$\|u(t) - v_3(t); L^2\| \rightarrow 0$$

as $t \rightarrow \pm\infty$.

We use the method of Hayashi and Naumkin [15,17]. The following ingredients are new and necessary to provide improvements, however. First, our method depends exclusively

on a contraction argument and is independent of a contradiction argument in [15,17,18]. Secondly, our method depends exclusively on the generators of Galilei transforms and is independent of the usual regularity argument. This enables us not to impose any regularity assumption on the Cauchy data. Thirdly, our argument treats the L^2 norm and homogeneously weighted norms separately for the Cauchy data as well as for solutions. This enable us not to impose smallness of the L^2 norm of the Cauchy data. For instance, data of the form $\varepsilon^{-1}\psi(\varepsilon^{-1}x, \varepsilon^{-1}y)$ with $\varepsilon > 0$ sufficiently small and $\psi \in L^2 \cap \mathcal{F}(\dot{H}^{\alpha,\beta})$ fall within the scope of Theorem 2.

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